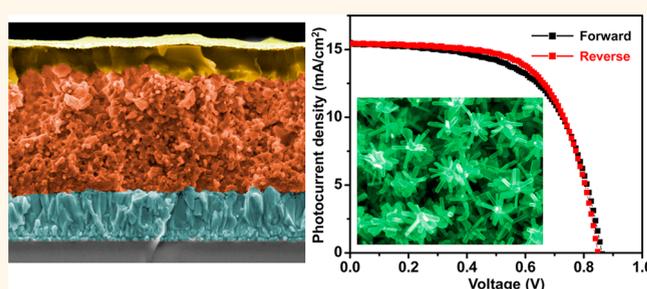


# Development of Lead Iodide Perovskite Solar Cells Using Three-Dimensional Titanium Dioxide Nanowire Architectures

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**ABSTRACT** Three-dimensional (3D) nanowire (NW) architectures are considered as superior electrode design for photovoltaic devices compared to NWs or nanoparticle systems in terms of improved large surface area and charge transport properties. In this paper, we report development of lead iodide perovskite solar cells based on a novel 3D TiO<sub>2</sub> NW architectures. The 3D TiO<sub>2</sub> nanostructure was synthesized via surface-reaction-limited pulsed chemical vapor deposition (SPCVD) technique that also implemented the Kirkendall effect for complete ZnO NW template conversion. It was found that the film thickness of 3D TiO<sub>2</sub> can significantly influence the photovoltaic performance. Short-circuit current increased with the TiO<sub>2</sub> length, while open-circuit voltage and fill factor decreased with the length. The highest power conversion efficiency (PCE) of 9.0% was achieved with ~600 nm long 3D TiO<sub>2</sub> NW structures. Compared to other 1D nanostructure arrays (TiO<sub>2</sub> nanotubes, TiO<sub>2</sub>-coated ZnO NWs and ZnO NWs), 3D TiO<sub>2</sub> NW architecture was able to achieve larger amounts of perovskite loading, enhanced light harvesting efficiency, and increased electron-transport property. Therefore, its PCE is 1.5, 2.3, and 2.8 times higher than those of TiO<sub>2</sub> nanotubes, TiO<sub>2</sub>-coated ZnO NWs, and ZnO NWs, respectively. The unique morphological advantages, together with the largely suppressed hysteresis effect, make 3D hierarchical TiO<sub>2</sub> a promising electrode selection in designing high-performance perovskite solar cells.



**KEYWORDS:** three-dimensional nanowires · titanium dioxide · CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> · perovskite · solar cells

Rapid charge transport and high surface area are preferable characteristics for solar energy harvesting including photoelectrochemical (PEC) cells and dye-sensitized solar cells (DSSCs).<sup>1</sup> Single crystalline nanowires (NWs) and mesoporous nanoparticles (NPs) are representative structures with the merits of fast electron conducting and large surface area, respectively, whereas intrinsic limitations such as low surface area of NWs and poor electron mobility of NPs pronouncedly compensate their aforementioned advantages. Recently, three-dimensional (3D) hierarchical nanostructures have emerged as promising building blocks for photoelectrode design because they offer high-speed pathways as well as considerable surface area for sufficient chemical reaction or dye loading simultaneously.<sup>2–4</sup> Heterogeneous structures

including Si/TiO<sub>2</sub>,<sup>5–7</sup> Si/ZnO,<sup>8–10</sup> and CuO/ZnO<sup>11</sup> have demonstrated remarkable enhancement of efficiency in PEC water splitting compared to bare NW electrodes. Concurrently, homogeneous TiO<sub>2</sub>-phase 3D nanostructures were found to be effective in DSSCs development and exhibited a fast electron-transport rate over randomly oriented NPs.<sup>12–19</sup> Substantial research efforts have notably boosted the power conversion efficiency (PCE) through fabrication of delicate hierarchical TiO<sub>2</sub> nanostructures. However, the low absorption coefficient of traditional organic dye limits further enhancement of 3D TiO<sub>2</sub>-based DSSCs.<sup>20,21</sup>

Methylammonium lead iodide (CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>) is rising as a new star in the photovoltaic (PV) field because of its favorable direct band gap (1.55 eV) for good light absorption (entire visible light absorption), small

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binding energy (about 0.030 eV) of excitons, high carrier mobility ( $7.5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  for electrons and  $12.5 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  for holes), and long charge diffusion length (ranging from 100 to 1000 nm).<sup>22–29</sup> PCEs of around 15% have been frequently reported on the basis of both planar and mesoscopic cell configurations.<sup>30–37</sup> Advantages of incorporating mesoporous NPs include facilitating charge collection and alleviating the hysteresis effect,<sup>28,38,39</sup> which is a severe issue that could over- or underestimate PCEs.<sup>28,40–42</sup> Besides mesoporous NPs, NWs have also been intensively investigated as the framework for supporting perovskite.<sup>20,21,43,44</sup> However, to date, there is still no practical study of 3D NW-based perovskite solar cells (SCs) although tremendous merits are intuitively expected. For instance, the large absorption coefficient of perovskite is highly desired for 3D NW-based SCs. The stem of the 3D structure could serve as a rapid charge-conducting channel to facilitate electron collection and possibly reduce the hysteresis effect in analogy with NPs. Compared to regular NW arrays, the 3D branched NW structure is able to provide a much higher volume density of nucleation sites for the deposition of perovskite and thus allows larger photoactive materials loading. Additionally, the entire treelike morphology could increase the optical pathways by light scattering.

In this paper, we report a 3D  $\text{TiO}_2$  NW-based lead iodide perovskite SCs. A surface-reaction-limited pulsed chemical vapor deposition (SPCVD) technique was employed to fabricate the 3D  $\text{TiO}_2$  architecture,<sup>45</sup> where the  $\text{CH}_3\text{NH}_3\text{PbI}_3$  perovskite was deposited using the sequential deposition method.<sup>30</sup> The highest PCE yielded by the 3D  $\text{TiO}_2$  structure was found to be 9.0%, which was substantially higher than perovskite SCs made from ZnO NWs,  $\text{TiO}_2$ -coated ZnO NWs, and  $\text{TiO}_2$  nanotubes (NT). Over 8% of PCEs have been achieved on the basis of 3D  $\text{TiO}_2$  with length variations ranging from  $\sim 600 \text{ nm}$  to  $\sim 1.4 \mu\text{m}$ . Furthermore, enhanced light absorption at the visible light region, increased electron mobility, and low hysteresis effects were observed from the 3D  $\text{TiO}_2$ -based perovskite SCs.

## RESULTS AND DISCUSSION

3D  $\text{TiO}_2$  NW architecture was synthesized inside a homemade atomic layer deposition (ALD) system using ZnO NW arrays as the template and  $\text{TiCl}_4$  and  $\text{H}_2\text{O}$  as the precursors.  $\text{TiO}_2$  NT was formed after 10 cycles of deposition through the cation-exchange reaction governed by the Kirkendall effect.<sup>45</sup> Subsequently, the SPCVD process dominated the growth, forming high-density single crystalline  $\text{TiO}_2$  nanorods (NRs) on the entire  $\text{TiO}_2$  NT stem (see the Methods section for details). Figure 1a shows scanning electron microscopy (SEM) images of as-received treelike 3D  $\text{TiO}_2$  NW structures (inset of Figure 1a). The interspace between individual  $\text{TiO}_2$  nanostructures was from 160 to 430 nm, which was within the range of the grain size

distribution of perovskite acquired from the sequential deposition approach, implying possible full  $\text{CH}_3\text{NH}_3\text{PbI}_3$  coverage of the open space. TEM characterizations on a branched NW demonstrate the dense and uniform  $\text{TiO}_2$  NR distribution along the entire trunk with a density of  $\sim 70$  NRs per  $\mu\text{m}$  (Figure 1b). The trunk is polycrystalline with a hollow center (Figure 1c), which manifests the Kirkendall effect induced by different ionic diffusion rates of the ZnO core and  $\text{TiO}_2$  shell during the vapor-phase deposition.<sup>45</sup> The shell thickness of the  $\text{TiO}_2$  NT core was found to be  $\sim 12 \text{ nm}$ . A high-resolution TEM (HRTEM) image reveals the high-quality lattice of single crystalline  $\text{TiO}_2$  NRs which do not show any line or planar defects. Lattice spacing measured from the HRTEM image was 0.48 and 0.35 nm, perfectly matching the (002) and (011) planes of anatase  $\text{TiO}_2$ , respectively.<sup>46,47</sup> Besides the dominated anatase phase, rutile  $\text{TiO}_2$  was also observed from the X-ray diffraction (XRD) pattern (Figure 1d), which was likely formed in the stem portion where foreign elements, such as Zn, were involved.<sup>45</sup>

Figure 2 shows cross-sectional SEM images of 3D  $\text{TiO}_2$  NW architecture (Figure 2a),  $\text{CH}_3\text{NH}_3\text{PbI}_3$  perovskite-coated  $\text{TiO}_2$  (Figure 2b), and the complete PV devices (Figure 2c). The film thicknesses of perovskite were governed by the length of 3D  $\text{TiO}_2$  NWs. Four different lengths were selected ( $\sim 600 \text{ nm}$ ,  $\sim 1.0 \mu\text{m}$ ,  $\sim 1.4 \mu\text{m}$ , and  $\sim 1.7 \mu\text{m}$ ) and marked as i, ii, iii, and iv, respectively. The length of  $\text{TiO}_2$  was controlled by the length of ZnO NW templates through adjustment of the synthesis period.<sup>48</sup> After sequential deposition of  $\text{PbI}_2$  and  $\text{CH}_3\text{NH}_3\text{I}$ ,  $\text{CH}_3\text{NH}_3\text{PbI}_3$  perovskite could densely cover the entire 3D  $\text{TiO}_2$  NW architecture with a film thickness up to  $\sim 1.4 \mu\text{m}$  (Figure 2b, i–iii), indicating a large loading amount of photoactive perovskite. Such an efficient perovskite loading benefited from the high density  $\text{TiO}_2$  branch that serves as the nucleation site for perovskite formation. For the shortest  $\text{TiO}_2$  scaffolds, large-sized perovskite crystals with sizes of 120 to 450 nm were also formed and nearly covered the entire top area. Longer  $\text{TiO}_2$  structures possess larger space in between, and fewer perovskite crystals were observed on the top (Figure S1, Supporting Information). When the  $\text{TiO}_2$  thickness increased to  $\sim 1.7 \mu\text{m}$ , an uncovered region around the top of  $\text{TiO}_2$  could be observed (marked by the dashed white box in Figures 2b, iv, and 2c, iv). XRD patterns represent pure tetragonal phase of  $\text{CH}_3\text{NH}_3\text{PbI}_3$  perovskite for samples i–iv (Figure S2, Supporting Information), showing a good agreement with previously reported perovskite on  $\text{TiO}_2$  scaffolds.<sup>30,38</sup>

Good coverage of perovskite on samples i–iii renders a typical layered structure of perovskite SCs after filling the small pores inside photoanode with hole-transporting materials (HTM) (Figure 2c, i–iii). The silver electrode and photoanode were completely

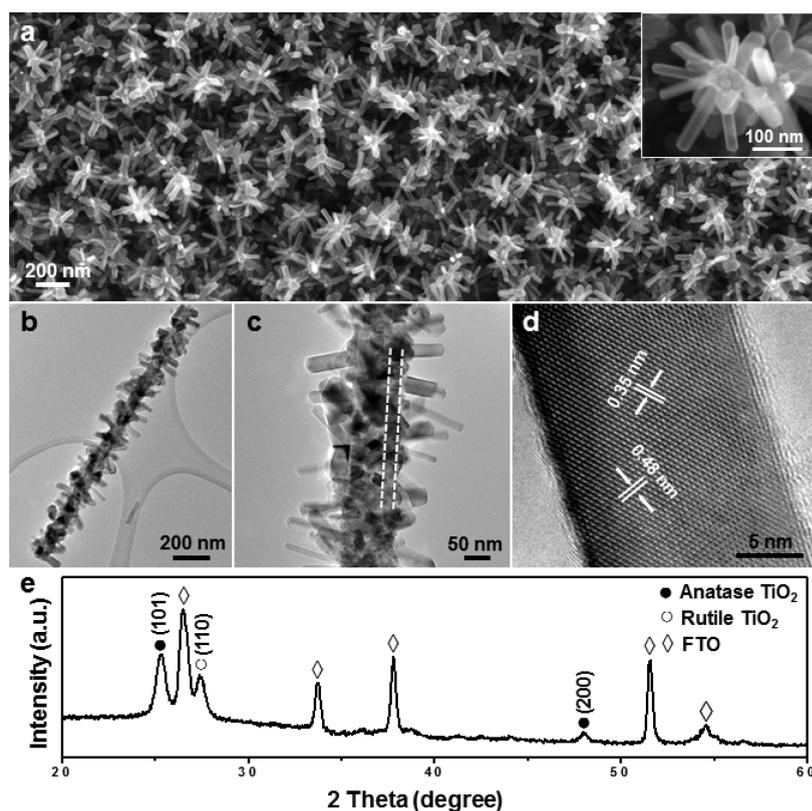


Figure 1. (a) Planar view of a 3D TiO<sub>2</sub> NW architecture. Inset is an enlarged SEM image showing the treelike branched NR structure. (b) TEM image of a branched TiO<sub>2</sub> NW showing dense and uniform coverage of TiO<sub>2</sub> NRs along the entire stem. (c) Higher magnification TEM image showing the single-crystalline NR branches and the polycrystalline NT trunk. (d) HRTEM image of a TiO<sub>2</sub> NR branch showing the perfect crystal lattice. (e) X-ray diffraction pattern of the 3D TiO<sub>2</sub> nanostructure revealing both anatase and rutile TiO<sub>2</sub> phases.

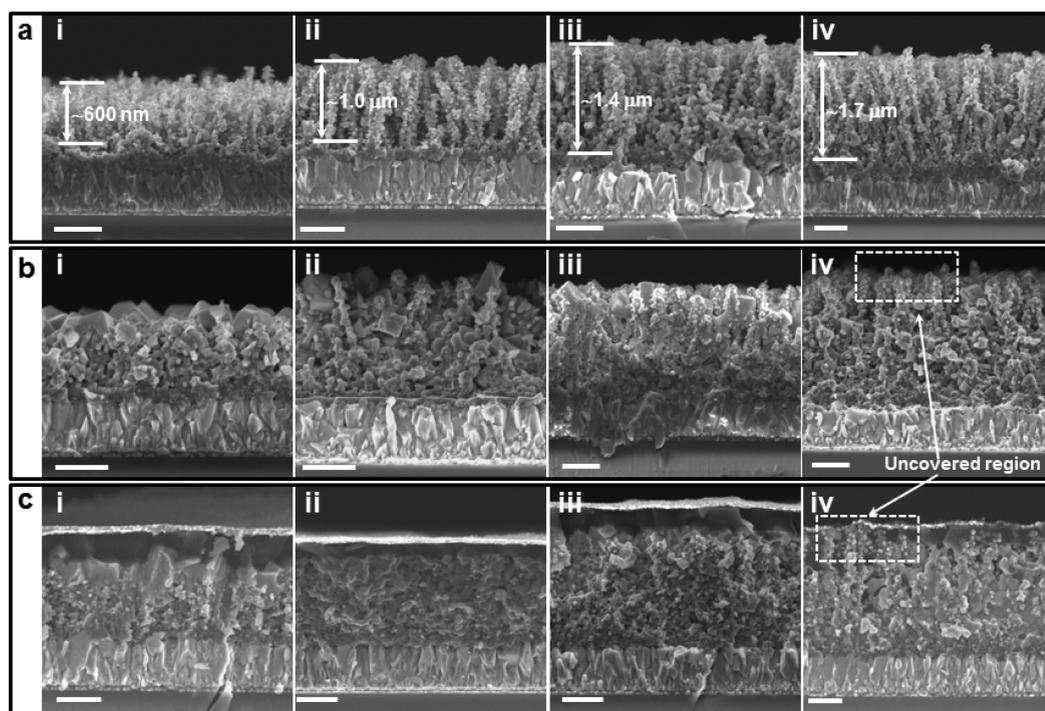


Figure 2. Cross-sectional SEM images of 3D TiO<sub>2</sub> nanostructures (a), CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> perovskite coated TiO<sub>2</sub> (b), and the final SC devices (c). The lengths of nanostructure are ~600 nm, ~1.0 μm, ~1.4 μm, and ~1.7 μm and are marked as i, ii, iii, and iv, respectively. Scale bars are 500 nm.

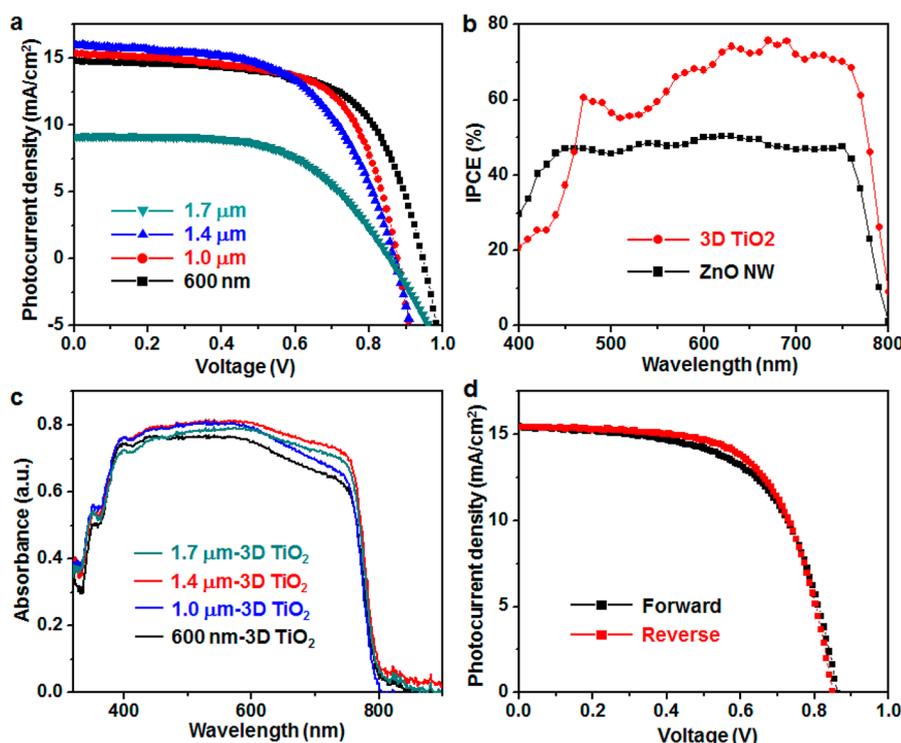


Figure 3. (a)  $J$ - $V$  curves of a series of 3D  $\text{TiO}_2$ -based perovskite SCs with different film thicknesses tested under 1 sun irradiation. (b) IPCE spectra of 3D  $\text{TiO}_2$  (red) and ZnO NW (black) based SCs with a film thickness of  $\sim 600$  nm. (c) UV-vis absorption spectra of  $\text{CH}_3\text{NH}_3\text{PbI}_3$  coated 3D  $\text{TiO}_2$  nanostructure with different film thicknesses. (d)  $J$ - $V$  curves of a cell made with  $\sim 1.4$   $\mu\text{m}$  3D  $\text{TiO}_2/\text{CH}_3\text{NH}_3\text{PbI}_3/\text{spiro-MeOTAD}/\text{Ag}$  measured by forward (black) and reverse (red) scans with 7.5 mV voltage steps and 200 ms delay time under 1 sun illumination.

TABLE 1. PV Performances of Perovskite SCs Fabricated Using 3D  $\text{TiO}_2$  NW Architectures with Different Lengths

$\text{TiO}_2$ length	best performance				avg of 20 cells			
	$J_{\text{sc}}$ ( $\text{mA}/\text{cm}^2$ )	$V_{\text{oc}}$ (V)	FF	$\eta$ (%)	$J_{\text{sc}}$ ( $\text{mA}/\text{cm}^2$ )	$V_{\text{oc}}$ (V)	FF	$\eta$ (%)
600 nm	14.8	0.95	0.64	9.0	$13.5 \pm 1.3$	$0.94 \pm 0.1$	$0.63 \pm 0.2$	$8.0 \pm 1.0$
1.0 $\mu\text{m}$	15.3	0.88	0.64	8.6	$13.9 \pm 1.4$	$0.87 \pm 0.1$	$0.61 \pm 0.3$	$7.4 \pm 1.3$
1.4 $\mu\text{m}$	16.0	0.87	0.58	8.1	$14.4 \pm 1.6$	$0.86 \pm 0.1$	$0.58 \pm 0.1$	$7.2 \pm 0.9$
1.7 $\mu\text{m}$	9.2	0.86	0.57	4.5	$7.3 \pm 1.9$	$0.81 \pm 0.5$	$0.49 \pm 0.8$	$3.1 \pm 1.4$

separated by the HTM layer. For sample iv, however, the open spaces are too large to be completely filled by HTM, and thus, a few NW tips touched the silver electrode (Figure S3, Supporting Information). The thickness of the HTM overlayer for samples i–iii was 289, 164, and 190 nm, respectively. The thickness of the HTM overlayer is inversely proportional to the open spaces that exist among the nanostructure scaffolds, and the largest thickness was obtained from the shortest 3D  $\text{TiO}_2$ . Similar overcoating results were also observed from mesoporous  $\text{TiO}_2$  NP<sup>49</sup> and ZnO NW-based<sup>44</sup> perovskite SCs.

The  $J$ - $V$  curves of perovskite SCs with different  $\text{TiO}_2$  lengths under one sun illumination are represented in Figure 3a. In general, PCE decreased with the increase of film thickness (Table 1). This trend has also been reported from other nanostructured perovskite SCs and was suggested to be a result of declining charge

generation efficiency, increased charge recombination, and raised charge-transport resistance.<sup>20,50</sup> The highest PCE of 9.0% was achieved from the SC made from  $\sim 600$  nm long 3D  $\text{TiO}_2$ , giving a short-circuit photocurrent density ( $J_{\text{sc}}$ ) of 14.8  $\text{mA}/\text{cm}^2$ , an open-circuit voltage ( $V_{\text{oc}}$ ) of 0.95 V, and a fill factor (FF) of 0.64. The average PCE of 8.0% was received by  $\sim 600$  nm long 3D  $\text{TiO}_2$  with a variation of 1.0% (Table 1). The IPCE spectrum (Figure 3b) exhibits large external quantum efficiency (EQE) of more than 70% across a wavelength range from 600 to 750 nm with a maximum of 75.7% attained at 670 nm. The onset value of 800 nm is consistent with the reported band gap of  $\text{CH}_3\text{NH}_3\text{PbI}_3$  (1.55 eV).<sup>23</sup> Integrating the product of AM 1.5G photon flux with IPCE spectrum yields a predicted  $J_{\text{sc}}$  of 14.5  $\text{mA}/\text{cm}^2$ , matching well with the measured  $J_{\text{sc}}$  of 14.8  $\text{mA}/\text{cm}^2$ .

It is important to note that the IPCE at the blue and green light region is relatively lower than typical

CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> SCs. This outcome is presumably related to the doped nature of the 3D TiO<sub>2</sub> nanostructure. Trace amounts of Zn residue have been observed in the TiO<sub>2</sub> trunk, which might enhance the electron conductivity and also narrow the band gap of TiO<sub>2</sub>.<sup>45</sup> This claim is evidenced by the UV–vis absorption spectra of 3D TiO<sub>2</sub>, where considerable absorption was recorded in the visible light region, especially at the region between 400 and 600 nm (Figure S4, Supporting Information). Such an overlap in the absorption spectra would lead to competitive light absorption between the TiO<sub>2</sub> framework and photoactive perovskite, resulting in lower IPCE of the SC device. In contrast, TiO<sub>2</sub> NTs and ZnO NWs exhibited the regular threshold of 400 nm, consistent with the band gaps of undoped TiO<sub>2</sub> and ZnO. To further confirm this hypothesis, the IPCE spectrum of a ZnO NW-based perovskite SC is measured. The spectrum exhibits a characteristic shape of CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> SCs with ~50% IPCE across almost the entire visible light region (400–750 nm). This observation suggests that the CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> could absorb and utilize the entire visible light if no competitive absorption existed.

The PV performances of all 3D TiO<sub>2</sub>-based SCs are summarized in Table 1.  $J_{sc}$  exhibits a relatively small dependence on the TiO<sub>2</sub> length and increased slightly from 14.8 mA/cm<sup>2</sup> for ~600 nm to 16.0 mA/cm<sup>2</sup> for ~1.4 μm. Fundamentally,  $J_{sc}$  and IPCE are determined by light-harvesting efficiency, electron and hole injection efficiency, and charge collection efficiency.<sup>20</sup> Electron injection efficiency is expected to be independent with length since the interfacial property between TiO<sub>2</sub> and CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> should be identical regardless of length variation. Hole injection efficiency is proportional to the pore-filling fraction (PFF), which is commonly estimated by the thickness of the HTM overlayer. The thinner HTM overlayer indicates larger PFF and more effective hole injection. As shown in Figure 2c, a thinner HTM layer is observed from longer TiO<sub>2</sub>, enabling larger PFF and hole injection efficiency and further improving  $J_{sc}$ . The light-harvesting efficiency is determined by the amount of CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> and the optical path of incoming light. Notably, the longer TiO<sub>2</sub> nanostructure possesses more perovskite loading and affords higher light absorption. This is supported by the UV–vis absorption spectra (Figure 3c), where perovskite film on the longer TiO<sub>2</sub> framework illustrates enhanced light absorption in yellow and red light region (except the 1.7 μm sample due to its incomplete perovskite loading). To conclude, as the length of TiO<sub>2</sub> increases, efficiency of light harvesting and hole injection are expected to increase, while electron injection efficiency remains the same. Therefore, it is reasonable to see larger  $J_{sc}$  for longer TiO<sub>2</sub>. However, the thicker scaffold also induces enlarged competitive absorption by TiO<sub>2</sub>, which compensates the photocurrent enhancement and reduces the variation of  $J_{sc}$ .

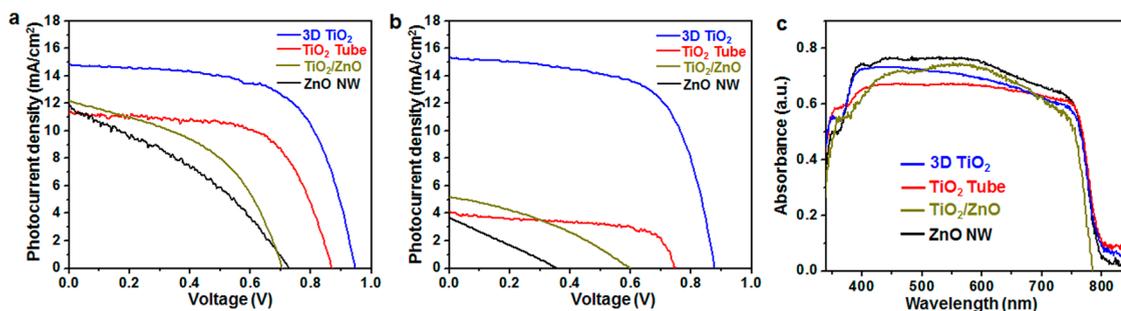
**TABLE 2. PV Performances of 3D TiO<sub>2</sub>-Based Perovskite SCs under Two Scan Directions**

scan direction	$J_{sc}$ (mA/cm <sup>2</sup> )	$V_{oc}$ (V)	FF	$\eta$ (%)
forward	15.5	0.86	0.60	8.0
reverse	15.4	0.86	0.63	8.3

$V_{oc}$  and FF are more significantly influenced by the TiO<sub>2</sub> length. Starting from 0.95 V of the ~600 nm TiO<sub>2</sub>,  $V_{oc}$  quickly declined to 0.88 V of the ~1 μm sample and then decreased to 0.86 V of the ~1.4 μm one. A similar trend of  $V_{oc}$  variation has also been reported in the case of mesoporous TiO<sub>2</sub> NPs<sup>50</sup> and TiO<sub>2</sub> NWs.<sup>20</sup> It was suggested that thicker photoanode with higher surface area would increase the recombination current and therefore lead to a lower  $V_{oc}$ .<sup>50</sup> FF gradually deteriorated from 0.64 to 0.57 with increased nanostructure thickness, which is a consequence of the lower  $V_{oc}$  and higher electron-transport resistance.<sup>20</sup>

The 3D TiO<sub>2</sub> nanostructures were found to be effective in minimizing the hysteresis effect in perovskite SCs. The hysteresis effect has recently been recognized as a significant issue that could over- or underestimate the PCE value. Mesoporous TiO<sub>2</sub> NPs was reported to be superior for suppressing the hysteresis compared to planar device configurations.<sup>28,38</sup> Similar to the NP systems, PV performance of 3D TiO<sub>2</sub>-based SCs demonstrated a very small hysteresis effect at a scanning delay time of 200 ms as shown in Figure 3d and Table 2.  $J$ – $V$  curves of the forward and reverse scans were almost identical. The  $J_{sc}$ ,  $V_{oc}$ , and FF under reverse scan are 15.4 mA/cm<sup>2</sup>, 0.86 V, and 0.63, respectively, yielding a PCE of 8.3%. Corresponding values for the forward scan are 15.5 mA/cm<sup>2</sup>, 0.86 V, and 0.60, respectively, yielding a PCE of 8.0%. The largely suppressed hysteresis effect indicates the strong reliability of measured PV performance. Conversely, planar perovskite SC showed a significant hysteresis effect with a PCE of 9.5% under reverse scan and a PCE of 4.4% under the forward scan (Figure S5 and Table S1, Supporting Information). For mesoporous TiO<sub>2</sub> NP based perovskite SCs, low hysteresis was only achieved with the less efficient device (PCE of 5.9%), and substantial variation of PCE under forward (6.9%) and reverse (10%) was observed for well-performed cells (Figures S6 and S7 and Table S1, Supporting Information).

To demonstrate the morphological advantage of the 3D TiO<sub>2</sub> architecture, bare TiO<sub>2</sub> NTs, ZnO NWs and TiO<sub>2</sub>-coated ZnO NWs were applied as scaffolds for fabricating perovskite SCs. Comparing the performance of these three nanostructured scaffolds clearly illustrated the morphological advantageous of the branched architecture in terms of perovskite loading, light scattering, and charge transport. Figure S8 (Supporting Information) shows the SEM images of single-crystalline ZnO NW arrays. Polycrystalline TiO<sub>2</sub> NTs were obtained by converting ZnO NWs following the



**Figure 4.**  $J$ – $V$  curves of perovskite SCs made from 3D  $\text{TiO}_2$  nanostructure,  $\text{TiO}_2$  NTs,  $\text{TiO}_2$ -coated ZnO NWs, and ZnO NWs with an identical film thickness of  $\sim 600$  nm (a) and  $\sim 1.0$   $\mu\text{m}$  (b). (c) UV–vis absorption spectra of  $\text{CH}_3\text{NH}_3\text{PbI}_3$ -coated 3D  $\text{TiO}_2$ ,  $\text{TiO}_2$  NT,  $\text{TiO}_2$ -coated ZnO NWs, and ZnO NWs with lengths of  $\sim 600$  nm.

cation-exchange-driven Kirkendall effect in SPCVD reactions.<sup>45</sup> The SEM image (Figure S9, Supporting Information) shows the one-dimensional shape of  $\text{TiO}_2$  NT with a rough surface. TEM characterizations in Figure S9 (Supporting Information) illustrate the hollow feature and thin polycrystalline shell. The planar and cross-sectional view of  $\text{CH}_3\text{NH}_3\text{PbI}_3$  perovskite coated ZnO NWs, and the final SC devices are shown in Figure S10 (Supporting Information). Three length scales of  $\sim 600$  nm,  $\sim 1.0$   $\mu\text{m}$ ,  $\sim 1.4$   $\mu\text{m}$  were investigated and marked as i, ii, and iii, respectively. The  $\text{TiO}_2$  NT-based film and device are shown in Figure S11 (Supporting Information). Figure S12 (Supporting Information) shows a TEM image of 29 nm thick  $\text{TiO}_2$ -coated ZnO NW and the corresponding perovskite film and final device configuration. As revealed by Figures S10 and S11 (Supporting Information), full coverage of  $\text{CH}_3\text{NH}_3\text{PbI}_3$  perovskite and the uniform layered device structure of ZnO NWs and  $\text{TiO}_2$  NTs were created only on the shortest nanostructures ( $\sim 600$  nm). A large amount of uncovered region (marked by dashed white box) can be clearly recognized when the length of nanostructure was  $\sim 1$   $\mu\text{m}$  and larger. This is strikingly different compared to the 3D  $\text{TiO}_2$  structure which can effectively support perovskite as thick as  $\sim 1.4$   $\mu\text{m}$ . Despite the coverage, pure phased  $\text{CH}_3\text{NH}_3\text{PbI}_3$  perovskite was received from both ZnO NW and  $\text{TiO}_2$  NT scaffolds (Figure S13, Supporting Information). Although reducing the space between NWs could also facilitate a perovskite loading ratio,<sup>44</sup> it may need to sacrifice the total volume of perovskite and bring challenges in NW synthesis, particularly for longer NWs (e.g.,  $>1$   $\mu\text{m}$ ).

Morphology-related PV performances were analyzed by comparing the  $J$ – $V$  curves of perovskite SCs fabricated from 3D  $\text{TiO}_2$ , ZnO NWs,  $\text{TiO}_2$  NTs, and  $\text{TiO}_2$ -coated ZnO NWs with the same length of  $\sim 600$  nm. 3D  $\text{TiO}_2$  achieved a PCE of 9.0%, resulting from a  $J_{\text{sc}}$  of 14.8  $\text{mA}/\text{cm}^2$ , a  $V_{\text{oc}}$  of 0.95 V, and a FF of 0.64. These values were all substantially higher than those of  $\text{TiO}_2$  NT and ZnO NW.  $\text{TiO}_2$  NTs exhibited a characteristic  $J$ – $V$  curve shape of typical  $\text{TiO}_2$ -based perovskite SCs, offering a reasonable FF of 0.63. With a  $J_{\text{sc}}$  of 11.3  $\text{mA}/\text{cm}^2$  and a

$V_{\text{oc}}$  of 0.87 V, a PCE of 6.2% were obtained. Nevertheless, the ZnO NW based device only reached a  $J_{\text{sc}}$  of 11.7  $\text{mA}/\text{cm}^2$ , a  $V_{\text{oc}}$  of 0.73 V, and a FF of 0.35, yielding a PCE of merely 3.0%. The attainable  $V_{\text{oc}}$  and FF of ZnO NW based device were apparently lower owing to the high rate of undesirable recombination compared with  $\text{TiO}_2$  photoanodes.<sup>51</sup> When the film thickness increased to  $\sim 1$   $\mu\text{m}$ , the PCE of hierarchical  $\text{TiO}_2$  remained as high as 8.6%. In contrast, performance of  $\text{TiO}_2$  NT and ZnO NW declined dramatically since no typical layered configuration can be maintained, leading to tremendous loss of photocurrent, photovoltage, FF, and PCE. SCs based on  $\text{TiO}_2$ -coated ZnO NWs yielded a PCE of 4.0% with a film thickness of  $\sim 600$  nm, and a similar deterioration of PCE was observed when its length increased to 1.0  $\mu\text{m}$  (Table S2, Supporting Information). The heterogeneous junction between the  $\text{TiO}_2$  film and ZnO NW may act as charge traps and lower the charge-transport efficiency.

The pronounced discrepancy of  $J_{\text{sc}}$  observed from 3D  $\text{TiO}_2$ , ZnO NWs, and  $\text{TiO}_2$  NTs at the film thickness of  $\sim 600$  nm are mainly attributed to the different light-harvesting efficiency and charge collection efficiency. As shown in Figure S1a (Supporting Information), there was an extra cap layer of  $\text{CH}_3\text{NH}_3\text{PbI}_3$  cuboids on the top of 3D  $\text{TiO}_2$ , which was found to be favorable in light harvesting, charge collection, and thus boosting efficiency.<sup>38,52</sup> Conversely, no over layer was observed on ZnO NWs (Figure S10a–i, Supporting Information) and  $\text{TiO}_2$  NTs (Figure S11a–i, Supporting Information), meaning lower perovskite loading and also a lack of the benefits from the bilayer structure. Moreover, branched architectures are believed to have higher light-scattering capability and, therefore, be advantageous in light harvesting.<sup>14–16</sup> The good light-harvesting property of 3D  $\text{TiO}_2$  was reinforced by the absorption spectra (Figure 4c), where 3D  $\text{TiO}_2$  clearly demonstrates larger light absorption over the entire visible light region.

The morphology-related performance variation can also be understood from the electronic transport properties and charge collection efficiency. Using conductive atomic force microscopy (C-AFM), we characterized

the conductivity of the 3D TiO<sub>2</sub>, TiO<sub>2</sub> NTs, and ZnO NWs from the tips of nanostructures to the bottom electrode. *I*–*V* measurements (Figure S14, Supporting Information) reveals that ZnO NW had the lowest resistivity ( $1.35 \times 10^{-4} \Omega \cdot \text{m}$  at a bias of 10 V) and polycrystalline TiO<sub>2</sub> NT demonstrated the highest resistivity ( $58.2 \Omega \cdot \text{m}$  at a bias of 10 V).<sup>51</sup> The high resistance is likely a result of the polycrystalline feature of the TiO<sub>2</sub> NTs. The resistivity of 3D TiO<sub>2</sub> ( $3.93 \times 10^{-2} \Omega \cdot \text{m}$  at a bias of 10 V) was higher than that of ZnO NW but was significantly lower than that of TiO<sub>2</sub> NTs. The substantially improved conductivity of the 3D TiO<sub>2</sub> is believed to be a result of the Zn impurity residing in the TiO<sub>2</sub> crystals during the cation-exchange processes. Therefore, in addition to the much larger surface area, the relatively high conductivity of the 3D TiO<sub>2</sub> architecture is also beneficial to charge separation and transportation, which is favorable for *J*<sub>sc</sub> and PCE improvements.

## CONCLUSION

In summary, CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> perovskite SCs were fabricated on the basis of 3D TiO<sub>2</sub> NW architecture. Sequentially deposited CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> perovskite was successfully integrated onto the 3D TiO<sub>2</sub> framework with a length variation from ~600 nm to ~1.4 μm. In general, by increasing the length of the 3D TiO<sub>2</sub> scaffolds, *J*<sub>sc</sub> was increased accordingly, while *V*<sub>oc</sub> and FF of cells were decreased. The enlarged *J*<sub>sc</sub> is attributed to the enhanced hole injection efficiency and increased light absorption from longer TiO<sub>2</sub> nanostructures. The length-related reduction of *V*<sub>oc</sub> and FF is

believed to be a result of decreased charge-generation efficiency, increased charge recombination, and higher electron-transport resistance. The best PCE was offered by the ~600 nm 3D TiO<sub>2</sub> scaffold with a value of 9.0%. Compared to 1D TiO<sub>2</sub> nanostructures made through the same process, the PCE of 3D TiO<sub>2</sub> was 1.5 times higher. It was even 2.3 and 2.8 times higher than that of TiO<sub>2</sub>-coated ZnO NWs and ZnO NW-based perovskite SCs, respectively. The high-density 3D TiO<sub>2</sub> architecture allowed higher perovskite loading and enhanced light harvesting compared with simple 1D nanostructure arrays. C-AFM also revealed a substantially improved conductivity compared to regular TiO<sub>2</sub> nanostructures, which promoted the charge-transport properties. In addition, PV cells comprising 3D TiO<sub>2</sub> are found to be favorable in suppressing the hysteresis effect that is often observed from planar perovskite SCs. Compared to the best performed CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> perovskite solar devices, the efficiency of the 3D TiO<sub>2</sub> based one is still a little lower. One possible reason is the competitive light absorption induced by the doped essence of 3D TiO<sub>2</sub>, which could negatively affect the light harvesting by perovskite and subsequently hamper *J*<sub>sc</sub> of the device. Nevertheless, the 3D hierarchical NWs morphology has convincingly shown advantages in terms of photoactive materials loading, electron delivering, and light scattering. It offers a promising candidate for understanding the property–performance relationship in perovskite solar cells and developing high-performance solar energy harvesting devices.

## METHODS

**Synthesis of ZnO NW Array Templates.** FTO glasses were cleaned by ultrasonication in a series of media including an alkaline, aqueous washing solution, acetone, and ethanol and eventually treated under oxygen plasma cleaning for 10 min. The dense TiO<sub>2</sub> blocking layer on the cleaned FTO substrate was made by spin coating the solution of titanium isopropoxide (700 μL), HCl (12 μL), and ethanol (10 mL) at 2000 rpm for 45 s, followed by annealing at 150 °C for 10 min and 500 °C for 30 min. Subsequently, a 50 nm thick ZnO film was coated by RF sputtering to act as seeds for ZnO NW growth. The substrate with seeds was immersed in the nutrient solution with the seeded surface facing downward. The nutrient solution comprised 25 mM zinc nitrate and 25 mM hexamethylenetetramine (HMT). ZnO NW length was controlled by the growing time, where 1 h and 10 min yielded ~600 nm NWs, 2 h for ~1.0 μm, 3 h and 20 min for ~1.4 μm, and 4 h for ~1.7 μm.

**Synthesis of 3D TiO<sub>2</sub> NW Architecture.** 3D TiO<sub>2</sub> NW architecture was synthesized in a homemade ALD system following the procedures described in our previous publication.<sup>45</sup> As-synthesized ZnO NW arrays were placed at the center of the ALD chamber. N<sub>2</sub> gas with a flow rate of 40 sccm was used to carry precursors into the chamber, generating a base pressure of ~4.1 Torr. TiCl<sub>4</sub> and H<sub>2</sub>O vapors were used as the precursors and pulsed into the furnace separately with a pulsing time of 500 ms and separated by 60 s N<sub>2</sub> purging. Therefore, one deposition cycle involves 500 ms of H<sub>2</sub>O pulse +60 s of N<sub>2</sub> purging +500 ms of TiCl<sub>4</sub> pulse +60 s of N<sub>2</sub> purging. The pressure change of TiCl<sub>4</sub> pulses was ~110 mTorr. The chamber temperature was

remained at 600 °C during the entire growth process. Subsequently, the furnace was cooled down naturally under N<sub>2</sub> flow. The hollow TiO<sub>2</sub> NT was created after 10 cycles of deposition. 400 cycles of deposition yielded the 3D TiO<sub>2</sub> NW architecture. TiO<sub>2</sub>-coated ZnO NWs were fabricated in the same ALD system by 400 cycles of TiO<sub>2</sub> deposition at 300 °C, yielding 29 nm thick polycrystalline TiO<sub>2</sub> film on ZnO NW surface.

**Perovskite Solar Cell Fabrication.** Sequential deposition approach was used to load perovskite onto 3D TiO<sub>2</sub>.<sup>30</sup> Methylammonium iodide (CH<sub>3</sub>NH<sub>3</sub>I) was first synthesized following the literature.<sup>53</sup> Specifically, methylamine (24 mL, 33 wt % in ethanol, Aldrich) reacted with hydroiodic acid (10 mL, 27 wt % in water) in ethanol (100 mL) under nitrogen atmosphere at room temperature for 2 h. Then, the solution was evaporated at 80 °C until white colored powder of CH<sub>3</sub>NH<sub>3</sub>I crystal was precipitated. PbI<sub>2</sub> (99%, Aldrich) was dissolved in *N,N*-dimethylformamide under stirring at 70 °C to yield a concentration of 460 mg/mL. The PbI<sub>2</sub> solution was infiltrated into the 3D TiO<sub>2</sub>, TiO<sub>2</sub> NT, ZnO NW, TiO<sub>2</sub>-coated ZnO NW, and mesoporous TiO<sub>2</sub> scaffolds by spin coating at 3000 rpm for 30 s and dried at 70 °C for 30 min. After being cooled to room temperature, the PbI<sub>2</sub>-coated samples were first prewetted in 2-propanol solution for 2 s and then immersed into solution of CH<sub>3</sub>NH<sub>3</sub>I in 2-propanol (10 mg/mL) for 70 s, followed by 2-propanol solution rinsing again. After drying in the air for several minutes, the films were heated at 70 °C for 30 min. HTM was then coated by spin coating at 2000 rpm for 45 s. HTM solution was prepared by dissolving 75 mg of (2,2',7,7'-tetrakis(*N,N*-di-*p*-methoxyphenylamine)-9,9-spirobifluorene)(spiro-MeOTAD), 28 μL of 4-*tert*-butylpyridine, and 17 μL of solution that contains 520 mg/mL of lithium

bis(trifluoromethylsulfonyl)imide in acetonitrile into 1 mL of chlorobenzene. Finally, 100 nm of silver was deposited by thermal evaporation on the top of the device.

A planar cell in Figure S5 (Supporting Information) was fabricated according to ref 33. Briefly, 0.2 g of  $\text{CH}_3\text{NH}_2$  and 0.578 g of  $\text{PbI}_2$  was mixed in 1 mL of anhydrous *N,N*-dimethylformamide by shaking at room temperature for 15 min, and a clear  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solution was obtained with a concentration of 45%. Sequentially, 100  $\mu\text{L}$  of  $\text{CH}_3\text{NH}_3\text{PbI}_3$  solution was dropped onto  $\text{TiO}_2$ -coated FTO substrate and spun at 5000 rpm. After spinning at 5000 rpm for 5 s, 100  $\mu\text{L}$  anhydrous chlorobenzene was quickly dropped onto the center of the substrate and then the color of substrate was changed to light brown. Total spin time was 30 s.

The mesoporous  $\text{TiO}_2$  layer composed of 20 nm-sized particles was deposited by spin coating at 2000 rpm for 60 s using a commercial  $\text{TiO}_2$  paste (Dyesol 18NRT, Dyesol) diluted in terpineol (1:3, weight ratio). After drying at 150 °C for 15 min, the  $\text{TiO}_2$  films were gradually heated to 500 °C and baked at this temperature for 30 min.

**Device Characterization.** *J*–*V* curves were measured with a solar simulator equipped with 150 W xenon lamp (6255, Newport) and Keithley 2636 source-meter. The light intensity was adjusted to 1 sun (100  $\text{mW}/\text{cm}^2$ ), and the active area was confined to 0.785  $\text{mm}^2$ . The IPCE was measured with a quartz tungsten halogen lamp modulated by a chopper wheel and a monochromator. The signal was recorded using a Stanford Research Systems lock-in amplifier at zero-bias and calibrated photodiodes (818 series, Newport). SEM characterizations were performed on Zeiss Leo 1530 field-emission microscope and TEM measurements were conducted on FEI TF30. X-ray diffraction pattern were acquired from the Bruker D8 Discovery with  $\text{Cu K}\alpha$  radiation.

**UV–vis Absorption Measurement.** Two machines were used to measure the absorption spectra. The absorption spectra on Figures 3c and 4c were recorded using an Evolution 220 UV–vis spectrophotometer with integrated sphere (ISA 220). Absorption was acquired by separately measuring the reflectance (%R) and transmittance (%T) spectra and using the formula  $A = 1 - \%R - \%T$ . A blank spectrum consisting of FTO only was measured and subtracted. The spectra of  $\text{TiO}_2$  and ZnO on Figure S4 (Supporting Information) were obtained via UV–vis spectroscopy (Agilent 8453) with a direct measurement of absorption.

**Conductive-Atomic Force Microscopy (C-AFM).** C-AFM data were recorded using XE-70 Park Systems. Samples were glued onto a steel disk. Part of the FTO surface was exposed and connected to the AFM stage using a Au wire. The AFM was operated in contact mode with a platinum cantilever. When the AFM tip touched the top of nanostructures, the *I*–*V* curve was collected by the AFM system. Voltage was swept from 0 to 10 V.

**Conflict of Interest:** The authors declare no competing financial interest.

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**Supporting Information Available:** Additional SEM images and XRD patterns of perovskite-coated 3D  $\text{TiO}_2$ , UV–vis absorption of  $\text{TiO}_2$  and ZnO scaffolds, *I*–*V* curves, IPCE and SEM images of planar and mesoporous  $\text{TiO}_2$ -based perovskite solar cells, SEM, XRD, and TEM characterizations of perovskite-coated ZnO NWs,  $\text{TiO}_2$  tube and  $\text{TiO}_2/\text{ZnO}$ , *I*–*V* measurements from C-AFM for comparing the conductivity of 3D  $\text{TiO}_2$ ,  $\text{TiO}_2$  tube and ZnO NWs, and tables for photovoltaic performance of control devices. This material is available free of charge via the Internet at <http://pubs.acs.org>

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